

1 Modelling autogenic morphodynamic processes in
2 meandering rivers with spatial width variations

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3 **Abstract.** Most morphodynamic models of river meandering assume spa-
4 tially constant width; depending on the intensity of spatial width variations,
5 different meandering styles actually exist, often associated with mid-channel
6 bars and islands. When intense enough, width oscillations characterize tran-
7 sitional planforms between meandering and braiding. We investigate, on a
8 modelling basis, morphodynamic feedbacks between spatial curvature and
9 width oscillations in river meanders, and related bedform patterns. Our re-
10 view of existing mathematical models suggests that width-curvature inter-
11 actions can be comprehensively analyzed by a hierarchy of models that de-
12 scend from a two-parameters perturbation solution of the governing depth-
13 averaged morphodynamic model. The focus is on in-stream, autogenic hydro-
14 morphodynamic processes, and not explicitly on bank processes. Curvature-
15 width interactions are fundamentally nonlinear: the perturbation approach
16 allows to investigate the key effects at the first nonlinear interaction. In me-
17 anders with initially constant width, curvature nonlinearly forces mid-channel
18 bar growth, promoting symmetrical bank erosion further downstream, pos-
19 sibly triggering width oscillations. These in turn can significantly affect the
20 process of bend stability and therefore condition the curvature dynamics. Wider-
21 at-bends meanders develop shorter bends and are morphologically more ac-
22 tive compared to equiwidth meanders, coherently with the few available field
23 observations. River evolution models aiming to separately simulate bank ero-
24 sion and accretion processes shall incorporate these autogenic flow-bed non-
25 linearities. Because of its focus on meandering morphologies close to the tran-

26 sition with braiding, the proposed approach can be taken as a novel, phys-
27 ically based viewpoint to the long-debated subject of channel pattern selec-
28 tion.

1. Introduction

29 Alluvial channel patterns are known to form a continuum between the poles of single-
30 thread meandering and multi-thread braiding (Figure 1). Although investigations on their
31 controlling factors date back to the nineteenth century (*Lokhtin* [1897]), the topic still rep-
32 resents one of the fundamental open issues in fluvial geomorphology and morphodynamics.
33 A huge amount of studies on river channel pattern and their environmental controls have
34 been proposed in the last fifty years. They differ in aim and approach and are either
35 qualitative or quantitative. Qualitative descriptive classifications (*Brice* [1982], *Mosley*
36 [1987]) distinguish morphological types on the basis of morphometric indicators. Process-
37 based qualitative classifications (*Schumm* [1977], *Church* [1992]), place more emphasis on
38 the main reasons for a river to develop the observed morphology. Quantitative channel
39 pattern predictors are of two major types. Discriminant functions separating conditions
40 of occurrence of straight, meandering, braiding and in some cases anastomosing rivers
41 have been both empirically (*Leopold and Wolman* [1957], *Ferguson* [1987] *Kleinhans and*
42 *van den Berg* [2011]) and rationally (e.g. *Huang et al.* [2004], *Eaton et al.* [2010]) derived.
43 Quantitative predictors based on migrating or steady bar theories in straight channels
44 (*Callander* [1969], *Engelund and Skovgaard* [1973], *Parker* [1976], *Fredsæ* [1978], *Crosato*
45 *and Mosselman* [2009]) predict channel pattern based on the linear solution of the depth-
46 averaged equations that govern the hydro-morphodynamics of open channel flow.

47 An overview of the existing approaches indicates at least three main areas with great
48 potential for scientific improvement. First, despite recent attempts at direct comparison
49 [*Kleinhans and van den Berg*, 2011], the different approaches have seldom been explic-

50 itly linked one another, and their integration is recognized to have great potential for
51 improved insight and predictive ability [*Eaton et al.*, 2010]. Second, most channel pat-
52 tern classifications and prediction models tend to sharply discriminate between different
53 river morphologies, although in practice the continuum of river patterns does not exhibit
54 sharp thresholds [*Ferguson*, 1987]. A broad spectrum of intermediate or transitional mor-
55 phologies (Figure 1b,c) indeed present typical features of both "pure" meandering, e.g.
56 sinuosity (Figure 1a) and braiding, e.g. the presence of islands and the locally associated
57 multi-thread tendency (Figure 1D). Third, the physics behind existing approaches is of-
58 ten relatively basic, particularly if compared with state-of-art morphodynamic modelling,
59 through which reproduction of many complex and nonlinear morphodynamic processes
60 has been achieved [*Seminara*, 2006].

61 The present paper focusses on river morphologies that can be considered intermedi-
62 ate between single-and multiple-thread, at the meandering-braiding transition, with the
63 primary aim to quantify key morphodynamic processes occurring in these transitional
64 patterns. It also presents a novel viewpoint on the channel pattern debate through
65 a physically-based, nonlinear modelling approach which is combined with a descriptive
66 channel pattern classification.

67 The paper builds on two basic assumptions. First, transitional morphologies can be
68 viewed as a tendency of single-thread meandering rivers to develop some degree of multi-
69 channel behavior. Therefore, insight in related morphodynamic processes can provide
70 complementary viewpoints to understand differences among channel patterns. Besides
71 being an intuitive concept, this is also suggested by the recent study of *Kleinhans and*
72 *van den Berg* [2011], who distinguish meandering rivers according to the dominant oc-

73 currence of "scroll bars" or "chute bars". Chute bars and chute cutoffs are associated
74 with more pronounced spatial width variations and tend to plot close to the meander-
75 ing/braiding discriminator curve at relatively high stream power (Figure 5 in *Kleinhans*
76 *and van den Berg* [2011]).

77 The second, related assumption is that the planform geometry of meandering and tran-
78 sitional river styles can be described through their spatial curvature and width distribu-
79 tions, with the amplitude of width oscillations increasing for transitional morphologies.
80 This emerges from descriptive channel pattern classifications. The one originally devel-
81 oped by *Brice* [1975] and lately refined by *Lagasse et al.* [2004] is particularly useful
82 for this purpose, because it essentially discriminates different styles of meandering based
83 on the intensity of spatial width oscillations. In Figure 2 we have reordered meandering
84 styles categories (A,...G) according to the absence / presence of spatial variations in chan-
85 nel width. It clearly appears that styles of alluvial meandering differ for the degree and
86 character of channel width and curvature variations, with transitional forms having the
87 highest degree of width oscillations. Figure 2 suggests that channel width may vary rather
88 systematically along a meander wavelength for *wider-at-bend* streams (classes B₂, C, D,
89 G₂) and more irregularly for class E streams. Therefore, the spatial width distribution
90 can be expected to play a key dynamic role in analogy with that of channel curvature in
91 equiwidth meanders.

92 A mechanistic understanding of the morphodynamics of transitional patterns requires
93 to investigate which physical processes control the presence of more or less pronounced
94 spatial variations in width. A nearly obvious consideration is that the width oscillations

95 are produced when point-bar growth and associated inner-bank advance cannot keep up
96 with outer-bank erosion [*Nanson and Hickin, 1983*].

97 A complementary and relatively unexplored argument is to which extent the spatial
98 oscillations in width may be significantly controlled by in-stream (or "autogenic") mor-
99 phodynamic processes, i.e. those controlling the bed-flow pattern in the central flow
100 region, where sidewalls boundary effects can be ignored. Field evidence has indeed been
101 provided [*Richards, 1976*] that systematic spatial variations of channel width can develop
102 in the absence of spatial variations in bank material composition and in close association
103 with central bars diverting the flow against the banks. This suggests that, besides near-
104 bank dynamics, autogenic hydromorphodynamic processes may significantly control the
105 planform and bedform evolution.

106 Moreover linear models based on autogenic, curvature-related hydromorphodynamic
107 processes in equiwidth meanders, and only very roughly accounting for bank processes,
108 have correctly reproduced key features of meander dynamics (e.g. *Ikeda et al. [1981]*,
109 *Blondeaux and Seminara [1985]*, *Struiksmma et al. [1985]*, *Odgaard [1986]*, *Howard [1992]*,
110 *Camporeale et al. [2007]*). Significant insight can therefore be expected from the inves-
111 tigation of autogenic processes associated with the contemporary presence of channel
112 curvature and width variations that describe the observed variety of meandering styles
113 and intermediate patterns.

114 The paper addresses this second argument and to this aim it focusses on the follow-
115 ing specific research questions. (1) Which physical processes are associated with tran-
116 sitional meandering styles? (2) Which cause-effect relationships are associated with the
117 contemporary presence of spatial curvature and width variations? (3) To which extent

118 morphological differences between equiwidth and transitional meandering styles can be
119 investigated through physics based instream morphodynamic models?

120 Answers to the above questions are explored first by reviewing the present state of art
121 about the role of curvature and width variations in single-thread rivers, with focus on
122 both field observations and mathematical modelling (Section 2). Some of these models
123 can be hierarchically derived within the same unified mathematical framework, with the
124 advantage that models at each order of approximation can be associated with a specific
125 physical process, thus fitting the purpose of the present review. The key parameters used
126 to quantify the degree and character of curvature and width variations are introduced in
127 Section 3 and quantified referring to field cases. Section 4 presents the unified hierarchical
128 framework of depth-averaged morphodynamic models for the flow-bed topography in a
129 meandering channel with spatial width variations. The solution approach is based on a
130 two-parameters perturbation expansion. Sections 5, 6 and 7 illustrate possible answers to
131 the specific research questions. Finally, results are harmonized and put in a broader per-
132 spective in the concluding Section, which also summarizes open issues for future research.

2. State of Art

133 Relatively few investigations have explicitly investigated the planform geometry, bed-
134 form characteristics and the morphodynamic processes typical of meandering channels
135 with spatial width variations. In the following the most relevant indications for the present
136 work, emerging from field observations and from mathematical modelling, are separately
137 reviewed.

2.1. Field Observations

138 Information about the causes for development and the dynamic role of width variations
139 in natural meandering rivers can be drawn from two major types of studies. First, analysis
140 of the large available meander bend dataset compiled by *Lagasse et al.* [2004] on the
141 basis of *Brice* [1982] original database provides quantitative evidence of morphodynamic
142 differences between meandering rivers that significantly differ in the degree of spatial width
143 oscillations. Second, a series of reach-scale field investigations (*Knighton* [1972], *Richards*
144 [1976], *Hooke* [1986], *Hooke and Yorke* [2011]) have documented key dynamic interactions
145 among mid-channel bars (bedforms) and channel width oscillations (planform) that have
146 occurred at the time scale of few channel-forming events in selected meandering river
147 reaches.

148 Data on hundreds of bends on 89 different meandering rivers in the U.S. are reported in
149 *Lagasse et al.* [2004]'s dataset, and are grouped according to the updated *Brice* meandering
150 river pattern classification (Figure 2). A relatively simple analysis of planform geometry
151 and migration properties of constant width meander bends (class B_1) and wider-at-bends
152 meandering reaches (class C) reveals that the observable geometrical discrepancies be-
153 tween the two classes might actually reflect differences in the controlling morphodynamic
154 processes. Wider-at-bends meandering reaches show a higher degree of morphological ac-
155 tivity with respect to those with constant width, as already argued by *Brice* [1982] (Fig-
156 ure 3a). For instance annual bend apex movement ζ^* exceeding 2% of the reach-averaged
157 channel width W^* can be measured in nearly 50 % of the analyzed wider-at-bends reaches
158 and only in 20% of the constant-width meander bends. The apex movement ζ^* has been

159 defined as in *Lagasse et al.* [2004]: a combination of extension, translation and expansion
160 of a meander loop.

161 Figure 3b (adapted from *Luchi et al.* [2011]) shows the cumulative frequency distribu-
162 tions of meander wavelength L^* scaled by the reach-averaged channel width W^* for the
163 two meandering styles C and B_1 . More than 60% of the constant width meanders' wave-
164 length exceed 15 times channel width, while only less than 40% of the wider-at-bends do.
165 As selection of different meander wavelengths is determined by different balances between
166 competing morphodynamic processes (e.g. *BollaPittaluga et al.* [2009]), this also suggests
167 that the two styles of meandering may reflect different autogenic physics.

168 More insight on the physics of meandering channels with spatial width variations is pro-
169 vided by field investigations carried out at the meander wavelength scale, complementing
170 the general tendencies illustrated in Figure 3. Width variations along meander bends
171 are often associated with mid-channel bars and vegetated islands (Figure 1b,c). The ob-
172 served morphology may represent a snapshot along a time sequence whereby the stream
173 may cut through and eventually abandon one of the low-flow branches [*Hooke and Yorke,*
174 2011]. Central bars and local widening not necessarily develop at bend apexes but can
175 be observed near inflection points (Figure 1C). Characteristic cycles of development of
176 mid-channel bars and longitudinal widening-narrowing sequences in single-thread streams
177 have been documented in the field by *Knighton* [1972], *Richards* [1976], *Hooke* [1986]
178 and *Hooke and Yorke* [2011] while other studies mainly focussed on single anabranches
179 of braided rivers (e.g. *Ashworth* [1996], *Klaassen et al.* [1993]), where the bed and banks
180 evolve at more comparable timescales [*Bertoldi and Tubino*, 2005].

181 *Knighton* [1972] described the time development of a mid-channel bar in the River Dean
182 (UK). The bar originally developed in a meander inflection region between two consec-
183 utive bends in correspondence of cross sectional overwidening related to bank erosion.
184 A two-channel pattern was observable at low flows, when the bar emerged and turned
185 into a temporary island (*Knighton* [1972], Figure 1). A sequence of competent flood
186 events caused a shift in the flow partition between the two branches, causing a temporary
187 multichannel behavior in what is normally classified as a meandering river.

188 A detailed description of the characteristic timescales and cycles of development of
189 mid-channel bars in a large number of bends of the River Dane (UK) has been provided
190 by *Hooke* [1986] and *Hooke and Yorke* [2011]. *Hooke* [1986] observed a typical sequence
191 whereby mid-channel bars started forming close to riffles, then related to cross-sectional
192 overwidening and after some formative events became attached to one of the banks. This
193 typical sequence has been lately confirmed by *Hooke and Yorke* [2011] in the same River
194 Dane (UK) having a characteristic timescale of few channel forming events (7-10 years).
195 Mid-channel bars position relative to bend apexes display a large scatter among different
196 bends, in most cases not coinciding with the bend apex.

197 The reach-scale observations of *Knighton* [1972] and *Hooke* [1986] confirm the detected
198 tendency of higher channel mobility and rapid bank erosion to occur in reaches with
199 more pronounced local widening. Also these observations confirm that the spatial width
200 oscillations are cyclic in time, reflecting a correlation with bed patterns oscillations. The
201 associated dynamics of mid channel bars contribute to the planform evolution of the most
202 active bends. Mid-channel bar growth followed by attachment to the inner bank promotes
203 bend growth due to the progressive abandonment of the inner branch [*Hooke*, 1986]. When

204 the opposite behavior applies, the flow concentrates into the inner channel and reproduces
205 an analogous dynamics to that of chute cutoffs occurring across braid bars in multi-thread
206 rivers [*Ferguson et al.*, 1992] and limits the increase of sinuosity. Although cross-sectional
207 overwidening has been identified as the primary mechanism responsible for the presence
208 and development of mid-channel bars (*Knighton* [1972], *Hooke* [1986]), however still it
209 is not clear whether bank erosion has a forcing or following function (*Lewin* [1981]) or
210 to which degree bed and banks interact dynamically through a mutual feedback process
211 (*Bertoldi and Tubino* [2005]).

212 Overall, the available field observations suggest significant dynamic differences between
213 equiwidth and transitional meandering morphologies, provide wavelength-scale description
214 of key physical processes and suggest that autogenic - or in-stream - processes occurring
215 in the central flow region can exert a significant control on the observed morphologies.
216 However a clear detection of cause-effect relationships between the competing mechanisms
217 and the establishment of a clearer legacy with meandering pattern descriptors requires a
218 comprehensive theoretical framework that still needs to be developed.

2.2. Mathematical Modelling

219 Several types of mathematical models suitable to investigate morphodynamic processes
220 in meandering rivers with spatial width variations have been developed in the past two
221 decades. The focus and approaches of these models differ significantly. A first distinction
222 can be made on the basis of the method used to compute the flow and bed topography
223 fields: both analytical and numerical solutions of the shallow water and sediment transport
224 equations have been proposed, as well as reduced complexity models. A second distinction
225 can be made between models incorporating in detail the physics of bank processes and

226 those incorporating them only indirectly through a linear, Partheniades-type relationship
227 between the lateral channel migration rate and the near-bank excess longitudinal velocity
228 or shear stress (*Partheniades and Paaswell* [1970], *Ikeda et al.* [1981]). Based on this, we
229 suggest to group existing models into four categories as reported in Table 1.

230 Since decades, most modelling-based physical insight on river meandering has been ob-
231 tained by means of simplified analytical solutions of morphodynamic models obtained
232 through perturbation methods [*Holmes*, 1995]. These approaches provide the foundation
233 for all models in categories 1 and 2 of Table 1. The well known bend theory originally
234 proposed by *Ikeda et al.* [1981], *Hasegawa* [1977], is obtained by linearizing the equa-
235 tions governing flow, sediment transport and the planform evolution. Among the major
236 achievements, the original bend theory and its subsequent refinements have allowed pre-
237 dicting the characteristic spatial scales of developing meanders [*Edwards and Smith*, 2002;
238 *Frascati and Lanzoni*, 2009], understanding conditions under which a meander can grow
239 towards mature loops [*Parker et al.*, 1983; *Blondeaux and Seminara*, 1985]; reproducing
240 typical meander loop migration rates [*Crosato*, 2009] and shapes [*Seminara et al.*, 2001;
241 *Zolezzi et al.*, 2009]; investigating the nature of meander instability [*Lanzoni and Semi-*
242 *nara*, 2006]; distinguishing stable and unstable bends in field cases [*Luchi et al.*, 2007];
243 and the dominant direction of upstream/downstream 2D morphodynamic influence [*Stru-*
244 *iksma et al.*, 1985; *Zolezzi and Seminara*, 2001; *Zolezzi et al.*, 2005]. When coupled with
245 planform evolution models at larger time scales, long-term simulation of of meander flood-
246 plain development [*Howard*, 1992; *Sun et al.*, 1996] and assessment of possible existence
247 of chaotic behaviors [*Stolum*, 1996; *Frascati and Lanzoni*, 2010] have been also achieved.

248 The rationale behind the application of perturbation methods is that flow and bed
249 topography deformation in curved channels can be assumed to be small with respect to the
250 flow velocity and water depth that would occur in a straight channel with flat bed fed by
251 the same discharge and with the same slope, width and sediment size. This is theoretically
252 justified because curvature is typically a small parameter - channel radius of curvature
253 is much larger than reach-averaged width - and planform geometry is slowly varying in
254 many meander bends. Most models used to study bend stability and meander planform
255 evolution are therefore linear [*Odgaard*, 1986; *Johannesson and Parker*, 1989; *Crosato*,
256 2008], the perturbation expansion being truncated at the first order of approximation. The
257 effect of flow nonlinearities has been studied by *Seminara and Tubino* [1992] who extended
258 the analysis at the third order of the perturbation expansion. *Seminara and Solari* [1998],
259 lately extended by *BollaPittaluga et al.* [2009] developed a slightly modified perturbation
260 approach that relaxed the assumption of small amplitude flow and bed perturbations thus
261 allowing a more complete treatment of nonlinear effects. *Tubino and Seminara* [1990]
262 modelled the conditions under which the migration of free bars is ceased in channels with
263 variable curvature.

264 Planform evolution models of river meandering relate a representative near-bank shear
265 stress or excess velocity to the lateral channel migration rate. Nearly all these models as-
266 sume the river width to keep constant in space and time, imposing that the bank retreat
267 rate is equal to the bank advance rate. This has been justified as a long term requirement
268 for meandering rivers and has received some support (*Pizzuto and Meckelnburg* [1989])
269 from field observations on rivers with fairly uniform cohesive banks. The assumption of
270 constant width has been relaxed only quite recently. Considering the significant scientific

271 improvement obtained through meander models with constant width, it can be expected
272 that allowing the presence of spatial width oscillations within a similar modelling frame-
273 work can lead to improved insight of the morphodynamics of transitional morphologies.

274 The linear analysis of *Repetto et al.* [2002] focussed on the steady flow-bed topogra-
275 phy deformation in straight channels with regular width oscillations, giving rise to mid-
276 channel bars and potential braiding initiation when planform instability occurs. *Repetto*
277 *and Tubino* [1999] investigated the nonlinear interaction between free migrating bars and
278 steady central bars in the same planform configuration. In the case of curved channels,
279 *Luchi et al.* [2010b] have modelled the linear and nonlinear development of mid-channel
280 bars in meandering channels, which are typical features of transitional morphologies be-
281 tween meandering and braiding; *Frascati and Lanzoni* [2011] have extended the theory
282 of *Repetto et al.* [2002] to channels with arbitrarily varying width, and applied this to
283 meandering streams; *Luchi et al.* [2011] have analyzed the dynamic effect of spatial width
284 oscillations on the process of meander bend stability; *Luchi et al.* [2012] have extended
285 the nonlinear model of *BollaPittaluga et al.* [2009] to meanders with spatial width os-
286 cillations. Linear models for equiwidth meandering channels [*Camporeale et al.*, 2007]
287 together with these more recent analytical theories for meanders with spatial width vari-
288 ations separately address the key physical processes emerging from the overview of field
289 investigations (Section 2.1).

290 The present review aims to show that these models can be derived within a unified,
291 hierarchical mathematical framework and that the resulting theoretical picture provides
292 a quantitative physical insight on transitional morphologies. Moreover, because these
293 models predict long term, "dynamic equilibrium" system tendencies, they are particularly

294 suitable to be conceptually linked with descriptive channel pattern classifications, with a
295 particular focus on the different styles of meandering. This will be presented in the next
296 sections, together with some unpublished data and novel applications, to attempt a first
297 answer to the research questions posed in the introduction.

298 Many other models can produce spatial variations in width, although the difference
299 in scales of application, methods used for the mathematical solution and specific focus
300 make them less suitable to be included within a hierarchical mathematical framework.
301 Nevertheless, highlighting their key properties is needed in the present review to provide
302 an overall modelling picture for transitional river morphologies.

303 Category (2) in Table 1 includes the approaches of *Chen and Duan* [2006], *Parker*
304 *et al.* [2011], *Motta et al.* [2012] and *Eke and Parker* [2011], which represent the most
305 promising and recent attempts to couple instream morphodynamics with a detail physical
306 description of the dynamics of bank regions. Their outcomes cannot be considered as fully
307 consolidated yet; therefore they will be discussed in the concluding section in relation to
308 the open research perspectives.

309 Widening in meandering channels has been also simulated by a series of numerical
310 models that couple two-dimensional, depth-averaged models of flow and bed topography
311 in movable computational grids with process-based bank erosion models. These have
312 been grouped under category (3) in Table 1 (e.g. *Mosselman* [1998], *Nagata et al.* [2000],
313 *Darby et al.* [2002] *Jang and Shimizu* [2005], *Rüther and Olsen* [2007]). The algorithms
314 have been written both in boundary-fitted non-orthogonal coordinate systems (following
315 *Mosselman* [1991]) and on unstructured meshes [*Rüther and Olsen*, 2007], who used a 3D
316 CFD approach. Models within this category can reproduce local widening in meandering

317 rivers and mainly differ for the level of detail at which the physical processes causing bank
318 erosion are incorporated and coupled with the flow-bed evolution model. For instance
319 *Duan and Julien* [2005] separate the calculation of bank erosion from the advance of bank
320 lines, using the parallel bank failure model for non-cohesive bank material and solving
321 the near-bank mass conservation equation. *Darby et al.* [2002] consider the deposition of
322 failed bank material at the toe of the bank and its subsequent removal. They show that
323 the lateral sediment input due to bank collapse might lead to the development of a wider,
324 shallower cross-section with respect to channels with fixed banks.

325 A common gap for the investigation of spatial width variations in meanders is the
326 absence of a sound formulation for bank accretion, which is essentially accounted for
327 as point bar deposition. Adequate modelling of bank accretion represents one of the
328 major unsolved issues in river morphodynamic modelling [*Mosselman*, 2011] and is almost
329 absent in all models belonging to the four categories of Table 1. An interesting attempt
330 to explicitly model bank accretion has been made by *Coulthard and Wiel* [2006] when
331 simulating the planform evolution of river meanders through the CAESAR model based
332 on cellular automata. Although the physics of bank erosion and accretion in *Coulthard*
333 *and Wiel* [2006] is strongly simplified - as for the flow and bed erosion/deposition model
334 - the two banks processes are decoupled and separately described. Lateral bank erosion is
335 computed as proportional to local channel curvature, derived on a cell by cell basis; bank
336 accretion is simulated through two slightly different relationships for the lateral sediment
337 exchange between the eroding and the accreting bank. Although the methods haven't been
338 tested in detail, they are both based on the critical role played by the lateral exchange of
339 sediments between the river and its floodplain for the spatial and temporal evolution of

340 channel curvature and width. Such exchange in meandering rivers is intrinsically related
341 to the process of spatially variable width adjustment: meander geometry itself implies
342 that more sediment is eroded at the outside edge than can be deposited on the inside, as
343 pointed out by *Lauer and Parker* [2008]. The meandering channel simulated by *Coulthard*
344 *and Wiel* [2006] is wider-at-bends, with sharper bends tending to become wider. To
345 the Authors' knowledge the work of *Coulthard and Wiel* [2006] is, strictly speaking, the
346 only example of reduced-complexity model applied to simulate the planform evolution
347 of river meandering with spatially variable width. In a broader sense, *Jagers* [2003],
348 building also on *Klaassen et al.* [1993], proposed an object-based simulation model -
349 "branches model" - for the planform evolution of braided streams, which also reproduced
350 the planform evolution of the weakly meandering anabranches of the braided network.
351 The two approaches of *Jagers* [2003] and of *Coulthard and Wiel* [2006] are grouped within
352 reduced-complexity models - category (4) - in Table 1.

3. Quantification of Width and Curvature Variations in Meanders

353 The modelling focus of the present paper requires a quantitative knowledge of the prop-
354 erties of the spatial series of curvature and width variations. A considerable amount of
355 studies has been devoted to the definition and analysis of the channel axis curvature and
356 of its longitudinal variations since decades (e.g. *Kinoshita* [1961], *Ferguson* [1973]) and
357 methods for computations based on aerial and remotely sensed images are continuously
358 being improved (*Güneralp and Rhoads* [2008]). The same does not apply to width varia-
359 tions in meandering rivers, on which very few quantitative studies have been concentrated
360 so far.

361 Most available data refer to reach-averaged values of channel width, for instance from
362 hydraulic geometry studies that provide bulk relationships but do not capture the spatial
363 fluctuations of width at the meander wavelength scale. The definition itself of meander
364 width is not obvious. Computation of width from aerial images can pose problems of
365 subjective interpretation that may prevent obtaining significant information in some cases.
366 Direct field estimates of bankfull conditions can also show some degree of subjectivity
367 where riparian areas are mainly grassland or covered with sparse plants especially close
368 to gently sloping convex banks. *Luchi et al.* [2010a] recently proposed a more objective
369 approach based on the application of non-uniform 1D steady flow model with fixed bed
370 that can be used when bed topography data are known with enough detail.

371 A look at the planform of free meandering streams suggests two useful assumptions for
372 morphodynamic modelling purposes. First, curvature and width are typically oscillat-
373 ing functions of the river arclength. Several reaches in alluvial river meanders typically
374 display periodic planform sequences whereby the channel axis can be locally described
375 by a sine-generated curve (*Langbein and Leopold* [1964]), a line whose curvature varies
376 sinusoidally as a function of the arclength. These forms can develop towards fattened and
377 skewed shapes or to compound loops (*Brice* [1974], *Hooke and Harvey* [1983]) that may be
378 reproduced through the inclusion of odd higher harmonics (*Kimoshita* [1961], *Zolezzi et al.*
379 [2009]) in the expression of the sine-generated curve. The styles of meandering reported
380 in *Lagasse et al.* [2004] classification (Figure 2) suggest that channel width might also
381 be described through a regularly oscillating function, with half the curvature wavelength,
382 namely for the "wider-at-bends" class.

383 At a first approximation, the planform of a meander with spatially variable width and
384 intrinsic wavelength L^* can therefore be represented by two oscillating functions of the
385 arclength describing the spatial distribution of channel width and curvature, with the
386 curvature wavelength being twice that of the width oscillations $L^* = 2L_w^*$. For sine-
387 generated meanders four main parameters are therefore needed: meander wavelength,
388 amplitude of curvature and width oscillations and the phase lag between the two functions.
389 To substantiate this assumption and to translate it into quantitative terms, it has been
390 tested against real meandering rivers data.

391 Besides the large *Lagasse et al.* [2004] dataset, an active meandering reach of the Rio
392 Beni in the Bolivian Amazon has been specifically chosen for this purpose. The Beni is a
393 large southern tropical river draining the Andean and sub-Andean ranges and flowing into
394 the Madeira River [*Gautier et al.*, 2007]. The reach is freely evolving, without significant
395 anthropic effects and with relatively homogeneous hydraulic and sediment conditions along
396 tenths of subsequent meander bends. It is of particular interest for the present study
397 because it shows wider-at-bends sections with frequent mid-channel bars and islands (see
398 also Figure 1b). Figure 4 shows the longitudinal distributions of the bankfull width,
399 computed assuming that areas of bare sediment at low flows are submerged at bankfull. In
400 most of the bends the width exhibits at least one significant peak between two subsequent
401 meander inflections. This agrees with analogous observations reported by *Luchi et al.*
402 [2011] based on data from a different reach of the same Rio Beni and is also consistent
403 with the outcomes of *Luchi et al.* [2010a] on a much different meandering river system
404 (River Bollin, UK).

The four parameters used to describe the river planform are expressed in dimensionless form for consistency with the modelling approach described in the next section. They are defined as follows:

$$\nu = \frac{W_0^*}{2R^*}; \quad \delta = \frac{(W_{max}^* - W_0^*)}{2(W_0^*)}; \quad \lambda = \frac{\pi W_0^*}{L^*}; \quad \lambda_w = 2\lambda. \quad (1)$$

In (1) a star (*) denotes dimensional quantities; λ represents the dimensionless meander wavenumber, taken as half the spatial frequency λ_w of width oscillations. Moreover L^* is meander intrinsic wavelength, W_0^* , W_{max}^* are the reach-averaged and maximum values of channel width, and ν , δ denote the dimensionless amplitudes of curvature and width oscillations, respectively. R^* is a typical measure of the channel axis radius of curvature; for a sine-generated meandering planform it has been often assumed equal to twice the minimum value at the bend apex. The spatial curvature $\mathcal{C}^*(s) = 1/R^*(s)$ and width distributions are made dimensionless with $2/W_0^*$ and W_0^* respectively. for a sine-generated meandering channel they can therefore be expressed through the dimensionless oscillating functions $\mathcal{C}(s)$ and $W(s)$ defined as follows:

$$\frac{W_0^* \mathcal{C}^*(s)}{2} = \mathcal{C}(s) = \nu [\exp(i\lambda s) + c.c] \quad (2)$$

$$\frac{W(s)}{W_0^*} = W(s) = 2 + 2\delta [\exp[i(\lambda_w s + 2\omega)] + c.c.] ; \quad (3)$$

with $c.c$ denoting the conjugate of a complex number and ω the phase lag between the widest and the most curved section.

The ratio $R^*/W_0^* = 1/(2\nu)$ has been subject of many quantitative analysis because of its fundamental relevance for meander dynamics. *Hickin and Nanson* [1984], *Hudson and Kesel* [2000], *Crosato* [2009], among others, pointed out the R^*/W_0^* range corresponding to the maximum meander migration rate, while *Hooke and Harvey* [1983] developed a

421 classification model for meander migration and loop shape based on the values of R^*/W_0^*
422 and of the meander path length ($L^* = \pi W_0^*/\lambda$).

423 Much less attention has been paid to quantify δ in meandering rivers. *Luchi et al.*
424 [2011] have estimated the values of δ by analyzing the standard deviation of the spatial
425 width distribution extracted for 31 bends of the Rio Beni. In most of the examined bends,
426 peak width values have been observed not to exceed 1.4 times the reach-averaged channel
427 width, corresponding to δ values (1) below 0.2.

428 We have tested these findings using data for wider-at-bends meanders (class C in Figure
429 2) reported in *Lagasse et al.* [2004] that contains information of maximum and reach-
430 averaged channel width. Equation (1) has been applied to all bends, which have been
431 grouped into classes of similar meander wavenumber λ . The outcomes, represented in
432 Figure 5 in the form of box plots, are consistent with the findings of *Luchi et al.* [2011]
433 on the Rio Beni. Average δ values are typically slightly lower than 0.1, corresponding to
434 the widest section less than 20% wider than the wavelength-averaged value. The largest
435 scatter is displayed by individual bends whose length is between 9 to 13 times the average
436 channel width ($0.25 < \lambda < 0.35$). The 84th percentile of all δ values never exceeds 0.13,
437 with outliers $\delta \simeq 0.2$ for very few bends (< 5%).

438 Figure 5b shows the same distributions for the dimensionless curvature amplitude ν
439 obtained from the same bends. Median values of ν show a regular increase with λ : this
440 might be associated with the tendency of channel curvature to vary more rapidly in space
441 when bends are shorter.

442 *Luchi et al.* [2011] have quantified also the distance between the widest section and the
443 bend apex. In the examined reach of the Rio Beni this relative distance can be relatively

444 high, normally keeping less than half the bend length (i.e. one quarter the meander
 445 wavelength L^*) for all bends. Only for a few bends is the maximum of channel width
 446 located closer to the meander inflections. Although the available data refer to only one
 447 river reach, no clear trend has been detected for the dimensionless phase lag ω between
 448 the width and curvature distributions.

449 Consistently with the few published data on the spatial series of width and curvature in
 450 real meandering rivers, the present analysis indicates that the dimensionless amplitudes of
 451 both distributions are "small". This suggests the suitability of a perturbation approach to
 452 seek for simplified analytical solution of the mathematical problem, which is formulated
 453 in the next section.

4. Hierarchy of Morphodynamic Models for Meandering Channels with Spatial Width Variations

The model formulated in the present section refers to the flow-bed topography deformation in meandering channels with spatially varying width. It essentially focusses on instream processes occurring in the central flow region. The dynamics of bank processes is accounted for through the simplified and widely adopted relation between bank migration and the near-bank instream flow field [*Partheniades and Paaswell, 1970*]. It stipulates that the bank erosion (accretion) rate is linearly related to the excess (defect) near-bank shear stress relative to a reach-averaged value. This assumption leads to the well known lateral channel migration law employed in most meander models (e.g. *Ikeda et al. [1981]*), which reads, in dimensionless form:

$$\zeta(s) = E [u_{lb}(s) - u_{rb}(s)], \quad (4)$$

454 where $\zeta(s)$ denotes the net rate of channel shift, assumed positive in the direction locally
 455 normal to the left bank and aiming outwards, s is the streamwise coordinate, u_{lb} and
 456 u_{rb} represent near-bank (left and right respectively) excess longitudinal velocity, and E is
 457 an empirical erosion coefficient which reflects bank properties [Hasegawa, 1989] compu-
 458 tational choices in the numerical scheme of meander planform evolution [Crosato, 2007]
 459 and floodplain heterogeneities [Güneralp and Rhoads, 2011].

4.1. Mathematical Formulation

460 The modelling tool chosen for the present work is based on a depth-averaged mor-
 461 phodynamic model forced by the spatial variations of channel width $W^*(s)$ and curvature
 462 $\mathcal{C}^*(s) = r_0^{*-1}$. This follows a rather established approach in meander morphodynamics and
 463 relaxes the assumption of a constant channel width typical of established meander evolu-
 464 tion models. The valley slope is assumed constant and the cohesionless bed is composed
 465 of uniform sediment size. The formulation refers to an intrinsic curvilinear coordinate
 466 system (s^*, n^*) that is right handed and orthogonal with the s^* axis locally downstream
 467 directed (Figure 6). The analysis is based on typical assumptions for large-scale river mor-
 468 phodynamic modelling. Namely it is referred to the central region of the cross section, it
 469 ignores the side boundary layers, it assumes a shallow water approximation and a slowly
 470 varying flow field in space. Moreover steady conditions for flow and bed topography are
 471 assumed considering a typical hierarchy of scales whereby planform geometry varies on a
 472 much longer time scale with respect to bed deformation, and to flow unsteadiness.

473 The depth averaged, steady formulation of the morphodynamic problem accounting
 474 for both curvature and width variations is formulated in dimensionless form with all the
 475 variables normalized through reach-averaged quantities in order to achieve more generality

476 and to facilitate comparison of the relative magnitude of different physical effects. The
 477 normalizing scales refer to a uniform flow in a straight channel that carries the same
 478 discharge, with the same mean sediment size d_s^* and slope S . The uniform flow depth is
 479 denoted with D_0^* .

The normalization procedure allows to point out the key dimensionless parameters: the aspect ratio β , the relative roughness d_s and the Shields stress τ_* . They are computed as follows using the reference uniform flow values, with Δ denoting the relative sediment submerged density:

$$\beta = \frac{W_0^*}{2D_0^*}, \quad d_s = \frac{d_s^*}{D_0^*}, \quad \tau_* = \frac{S}{\Delta d_s}. \quad (5)$$

480 The morphodynamic model can be expressed in the following form:

$$UU_{,s} + VU_{,n} + H_{,s} + \frac{\beta \tau_s}{D} + f_\alpha = f; \quad (6)$$

$$UV_{,s} + VV_{,n} + H_{,n} + \frac{\beta \tau_n}{D} + g_\alpha = g; \quad (7)$$

$$(DU)_{,s} + (DV)_{,n} = m; \quad (8)$$

$$Q_{s,s} + Q_{n,n} = p. \quad (9)$$

481 In equations (6, . . . , 9), U and V denote depth averaged flow velocity in the longitudinal
 482 and transverse direction respectively, H is the free surface elevation and D the local depth.
 483 (τ_s, τ_n) and (Q_s, Q_n) are the bottom shear stress and sediment rate vectors. Moreover f_α
 484 and g_α are homogeneous terms accounting for the effect of streamline curvature on the
 485 parameterization of secondary flows (see *Luchi et al.* [2011]), which vanish when secondary
 486 flows are parameterized using the channel axis curvature.

The forcing terms $\mathcal{F} = (f, g, m, p)$ in the right hand side of the differential system (6, ..., 9) can be written according to the same general structure:

$$\mathcal{F} = \nu \mathcal{F}_{10} + \delta \mathcal{F}_{01} + \nu^2 \mathcal{F}_{20} + \nu \delta \mathcal{F}_{11}, \quad (10)$$

487 which easily allows to distinguish between the forcing effects of width and curvature
 488 variations on the morphodynamic response of the system. The forcing effect of curvature
 489 is due to both first ($\mathcal{O}(\nu)$) and second ($\mathcal{O}(\nu^2)$) order terms and would appear also in
 490 meandering channels with constant width. Channel width variations force the system
 491 in the form of a first-order contribution $\mathcal{O}(\delta)$ that coincides with that corresponding to
 492 a straight channel with variable width. Moreover the $\mathcal{O}(\nu\delta)$ term represents the mixed
 493 effect due to width and curvature variations. Higher order terms in ν and δ and in their
 494 combinations are supposed to play a minor role. The complete expression of these forcing
 495 terms are reported in *Luchi et al.* [2011].

496 Finally, lateral boundary conditions to be associated with equations (6, ..., 9) impose
 497 that the lateral walls must be impermeable both to fluid and sediments.

4.2. Two-parameter Perturbation Solution Scheme

498 The data analysis of Section 3 has indicated that for wider-at bends meandering streams
 499 the dimensionless amplitude δ of width oscillations is a "small" number in analogy with
 500 the amplitude ν of curvature variations (Figure 5). Together with the structure of the
 501 forcing term (10), this suggests to employ a perturbation approach to study the dynamic
 502 role of width variations in meanders, analogous to that used in equiwidth meanders where
 503 the sole forcing effect is a spatially variable channel curvature. This slightly complicates
 504 the mathematical solution with respect to the equiwidth case because a two-parameters (ν

505 and δ) perturbation approach is required. In order to focus on the key physical processes
 506 and on the basis of the analysis described in the previous section, the solution is referred
 507 to an idealized meandering planform consisting of an indefinite sequence of sine-generated
 508 meander bends (*Langbein and Leopold [1964]*) and width oscillating with half the curvature
 509 wavelength (2,3).

510 The model (6 . . . 9) with boundary conditions is solved expanding the unknowns vector
 511 \mathbf{V} in powers of the two perturbation parameters ν and δ as follows:

$$\begin{aligned} \mathbf{V} &= \mathbf{V}_0 + \nu\mathbf{V}_{10} + \nu^2\mathbf{V}_{20} + \delta\mathbf{V}_{01} + \nu\delta\mathbf{V}_{11}; \\ \mathbf{V} &= (U, V, H, D); \quad \mathbf{V}_0 = (1, 0, H_0(s), 1); \\ \mathbf{V}_{kj} &= (U_{kj}, \dots, D_{kj}); (k = 0, 1, 2; j = 0, 1), \end{aligned} \tag{11}$$

512 a structure which is suggested by the right-hand side of (6, . . . , 9). On substituting (10)
 513 into (6, . . . , 9) and into the boundary conditions, linear differential systems are obtained
 514 at each order of approximation, which admit for exact analytical solutions, each one
 515 corresponding to a different physical mechanism:

- 516 • $\mathcal{O}(\nu)$: the flow and bed topography component linearly forced by curvature in me-
 517 anders with constant width;
- 518 • $\mathcal{O}(\nu^2)$: the second-order non linear component of flow and bed topography forced by
 519 curvature in meanders with constant width;
- 520 • $\mathcal{O}(\delta)$: the flow and bed topography component linearly forced by width variations in
 521 straight channels with variable width;
- 522 • $\mathcal{O}(\nu\delta)$: the first nonlinear interaction which expresses the mixed response for flow
 523 and bed topography in meanders with variable width.

524 Solutions at each of the above orders of approximation can be obtained separately
 525 because (ν, δ) can be assumed of the same order of magnitude and the spatial structure of
 526 the solutions with the same orders of magnitude ($\mathcal{O}(\nu)$, $\mathcal{O}(\delta)$ and $\mathcal{O}(\nu^2)$, $\mathcal{O}(\nu\delta)$) is different
 527 in both the (s, n) directions.

528 The next two sections have separate focuses on linear and nonlinear responses. The
 529 outcomes of linear solutions are quite established in meander modelling; combining these
 530 with the analysis of the nonlinear responses provides key ingredients to answer the research
 531 questions posed in the Introduction.

5. Linear Models

5.1. Equiwidth Meander: $\mathcal{O}(\nu)$

532 The bed deformation and lateral migration processes occurring in equiwidth meanders
 533 have been extensively modelled through linear models [*Camporeale et al.*, 2007]) commonly
 534 used to compute local migration rates in numerical models of long term planimetric evolu-
 535 tion (*Howard* [1992], *Seminara et al.* [2001], *Camporeale et al.* [2005]). Channel curvature
 536 is typically associated with laterally antisymmetric bed patterns: sequences of scour - de-
 537 position zones alternately spaced across the lateral and longitudinal direction. This causes
 538 an excess longitudinal velocity at one bank with respect to the reference uniform flow,
 539 and a symmetrical defect at the opposite bank. In terms of bedform, the $\mathcal{O}(\nu)$ solution
 540 reproduces the classical point bar morphology (see the corresponding sketch in the left
 541 column of Figure 12). It forces a laterally symmetrical pattern of the transverse velocity
 542 V_{10} and a lateral antisymmetric structure of the longitudinal velocity U_{10} and of the bed
 543 profile $\eta_{10} = F_0^2 H_{10} - D_{10}$ with the same longitudinal periodicity of the curvature (eqn.

544 2). The bed pattern is given by:

$$\begin{aligned}
 \eta_{10} &= \left[\alpha n + \sum_{j=1}^2 \gamma_j \sinh(\lambda_{j\nu} n) \right] e_1(s) + c.c. = \\
 &= \eta_{1\nu}(n) e_1(s) + \bar{\eta}_{1\nu}(n) \bar{e}_1(s),
 \end{aligned}
 \tag{12}$$

545 where $e_k = \exp(k i \lambda s)$ (k integer), $\alpha, \gamma_1, \gamma_2, \lambda_{1\nu}, \lambda_{2\nu}$ are complex numbers and an overbar
 546 denote complex conjugates. Mathematically the solution procedure is formally identical to
 547 that of *Blondeaux and Seminara* [1985], with some differences in the governing equations.

548 Meander evolution has been studied as a bend instability process considering that per-
 549 turbations of channel axis alignment with respect to a straight configuration may grow
 550 in time thus developing a meandering pattern. *Unstable* bends are defined as those that
 551 tend to further grow eventually leading to meander amplification. The stability of the
 552 small-amplitude sinusoidal perturbations of the channel axis depend on the phase lag be-
 553 tween the perturbation of the longitudinal velocity $U_{10}(s, n = 1)$ (see eqn. 4) and the
 554 curvature distribution. The highest migration rate is expected at the cross-section where
 555 the maximum difference between right and left bank excess longitudinal velocity occurs.

5.2. Straight Channels with Spatial Width Oscillations: $\mathcal{O}(\delta)$

556 Differently from curvature, spatial width variations force a laterally symmetrical flow-
 557 bed topography pattern. In this sense they represent a complementary planform effect
 558 with respect to curvature. As a result a central bar pattern is obtained, which tends more
 559 often, although not invariably, to appear in the widest sections, thus eventually promoting
 560 the growth of the width oscillations through a planform stability mechanism analogous to
 561 that of bend stability. *Repetto et al.* [2002] proposed an analytical model for the occurrence
 562 of mid-channel bars in straight channels with sinusoidal width variations and discussed the

563 role of central bars in triggering bifurcation of the stream and thus potentially initiating
 564 a braided pattern. The solution displays an antisymmetrical lateral structure of the
 565 transversal velocity V_{01} and a laterally symmetrical pattern of the longitudinal velocity
 566 U_{01} and of the bed profile $\eta_{01} = F_0^2 H_{01} - D_{01}$:

$$\begin{aligned} \eta_{01} &= [\alpha_1 \cosh(\lambda_{1\delta} n) + \alpha_2 \cosh(\lambda_{2\delta} n)] e_2 + c.c. = \\ &= \eta_{1\delta}(n) e_2(s) + \bar{\eta}_{1\delta}(n) \bar{e}_2(s). \end{aligned} \quad (13)$$

567 The longitudinal periodicity is that of the width variations (eqn. 3); $\alpha_1, \alpha_2, \lambda_{1\delta}, \lambda_{2\delta}$ are
 568 complex numbers. The bedform pattern at the $\mathcal{O}(\delta)$ is illustrated by the corresponding
 569 planform in the left column of Figure 12.

570 The concept of planform stability applied to this case implies that unstable planforms at
 571 the $\mathcal{O}(\delta)$ will tend to widen the widest section and thus can be interpreted as initiators of
 572 a multi-thread pattern. The analysis of *Repetto et al.* [2002] indicates that this planform
 573 instability theoretically occurs for long wavelengths ($\lambda_w \leq 0.15$, see also their Figure 22).

574 *Luchi et al.* [2010b] proposed that the linear mechanism theoretically studied by *Repetto*
 575 *et al.* [2002] can be a possible cause of mid-channel bar growth also when the channel axis
 576 has a meandering planform. This can occur when spatial width variations are already
 577 present, as when they originate because the advance rate of one bank cannot keep the
 578 pace of the opposite eroding bank (Figure 7). This can be referred to as a *width-forced*
 579 mechanism for mid-channel bar generation. It requires that different processes operate at
 580 opposite banks in the same cross-section. Such process has been documented by *Hooke*
 581 [1986] and *Hooke and Yorke* [2011], who indicated that widening preceded central bed
 582 deposition at most of the field sites.

583 Recently, the model of *Repetto et al.* [2002] has been extended to the case of arbitrarily
584 varying spatial width variations by *Frascati and Lanzoni* [2011]. This analysis is relevant
585 for meandering rivers because it allows to compute their complete linear response to
586 arbitrarily varying curvature and width: it results from truncating the expansion (10) at
587 the first two terms and relaxing the periodicity assumption of (2) and (3).

6. Nonlinear Models

6.1. Mid-channel Bars in Equiwidth Meanders: $\mathcal{O}(\nu^2)$

588 The presence of width variations and mid-channel bars in meanders is a typical chicken-
589 egg question: which of them does originate first? The linear, width forced mechanism
590 proposed by *Repetto et al.* [2002] and *Luchi et al.* [2010b] requires width variations to exist
591 prior to mid channel bars, like when opposite banks respectively experience accretion and
592 erosion.

593 A complementary process is the generation of width oscillation because of flow diver-
594 gence around existing mid-channel bars in originally equiwidth channels. This implies that
595 both banks in the same cross section are both subject to erosion, as in the section of the
596 River Bollin reported in Figure 9a. This has nearly vertical sidewalls and a transversally
597 symmetrical topographic high corresponding to a mid-channel bar pattern. Mathemati-
598 cally channel curvature can promote laterally symmetrical patterns at the $\mathcal{O}(\nu^2)$ solution,
599 which applies to a meander with constant width. Such possibility has been experimentally
600 substantiated by *Colombini et al.* [1992] on a sinusoidal, equiwidth meandering flume. The
601 analysis of the steady bed topography (see Figure 23.9 in *Colombini et al.* [1992]) pointed
602 out that the amplitude of the mid-channel bar component of the bed topography can be
603 up to half that of the point bar. Mid-channel bars are therefore not exclusively associated

with pre-existing spatial width variations, but can develop in equiwidth meanders because of nonlinear effects.

The longitudinal and lateral structure of the second order solution \mathbf{V}_{20} is determined by flow nonlinearities of the governing differential system (6 ... 9). At the $\mathcal{O}(\nu^2)$, as well as at each order of approximation, the solving differential system is linear. Therefore the structure of the solution closely matches that of the forcing terms. For instance, the term $UV_{,n}$ in the longitudinal momentum equation produces a forcing term due to the product of two first-order contributions:

$$\begin{aligned} V_{10}U_{10,n} &\Rightarrow (V_{1\nu}e_1 + \bar{V}_{1\nu}\bar{e}_1) (U_{1\nu,n}e_1 + \bar{U}_{1\nu,n}\bar{e}_1) \\ &= V_{1\nu}U_{1\nu,n}e_2 + V_{1\nu}\bar{U}_{1\nu,n}e_0 + c.c.; \end{aligned} \quad (14)$$

where we have employed analogous notations of equation (12) to express the lateral structures of V_{10} and U_{10} . In general, this implies that forcing terms for the longitudinal momentum, flow and sediment continuity equations are laterally symmetric. This is reflected in the solution for \mathbf{V}_{20} : in the s -direction the solution is the sum of one longitudinally oscillating component ($\propto e_2$) with twice the meander wavenumber (2λ) and one longitudinally invariant response ($\propto e_0$). The bed profile at the $\mathcal{O}(\nu^2)$ can be written as:

$$\eta_{20} = \left(\eta_{20}^{(2)}(n)e_2 + c.c. \right) + \eta_{20}^{(0)}(n)e_0; \quad (15)$$

being $\eta_{20}^{(2)}$ and $\eta_{20}^{(0)}$ even functions of n . The corresponding pattern is illustrated in the left column of Figure 12.

In order to relate the $\mathcal{O}(\nu^2)$ solution to processes occurring in a real meandering river, we have applied it to a reach of the River Bollin (NW England; Figure 8a), a gravel-bed meandering river whose high lateral mobility makes it particularly suitable to study me-

617 ander processes [Hooke, 2004]. The development of a mid-channel bar has been observed
618 between years 2008 and 2009 in several inflection regions like those appearing in Figure
619 8a. The subsequent bar growth can be assessed by visually comparing Figures 8b and
620 D, which have been taken before and after several channel-forming events. At forma-
621 tive conditions the Bollin has a nearly constant active channel width in space, i.e. the
622 width of the cross-section that is actually capable to transport sediments [Luchi *et al.*,
623 2012]. Therefore the observed mid-channel bar cannot be attributed to a width-forced
624 linear mechanism like that modelled by Repetto *et al.* [2002], but rather to a mechanism
625 operating in equiwidth meanders, like that corresponding to the $\mathcal{O}(\nu^2)$ solution.

626 Based on input parameters representative of formative conditions, the analytical solu-
627 tion easily allows to compute the position and the amplitude of the mid-channel bar and
628 related flow field along the meander according to such *curvature-driven* mechanism. In
629 Figure 8C the quantity φ_{η_2} is the phase lag between the most curved cross-section and the
630 peak of $\eta_{20}^{(2)}(n = 0)$, which corresponds to the mid-channel bar component of bed elevation
631 evaluated at the centerline. The continuous line indicates how the position of the central
632 bar top varies along half the meander wavelength (sketched in the left panel), depending
633 on the value of the dimensionless meander wavenumber λ . The wavenumber of bends in
634 the analyzed reach of the Bollin varies between 0.17 and 0.29, a range where the model
635 indicates a tendency to develop a central response close to meander inflection. This is
636 in good agreement with field analysis on the bed topography of Luchi *et al.* [2010a] that
637 pointed out the presence of an incipient mid channel bar at an inflection area between
638 two bends of similar wavenumber nearly equal to 0.19. The associated sections present an

639 evident laterally symmetrical, central bed topography pattern (Figure 9a) and the plot in
640 Figure 8D provides evidence of its permanence after some channel forming events.

641 An option to distinguish between the curvature-driven and width-forced mechanisms
642 for mid-channel bars with the aid of models can be based on the predicted location of
643 the mid-channel bar by the two mechanisms, as also illustrated by *Luchi et al.* [2010b]
644 and *Zolezzi et al.* [2011] referring to the meandering gravel-bed rivers Dane and Dean in
645 NW England. This has received some support also by the recent field study of *Hooke and*
646 *Yorke* [2011], which indicates that in the Dane mid-channel bars are width induced as
647 theoretically predicted.

648 Examining the location of the longitudinal velocity peak at the banks as predicted by
649 the $\mathcal{O}(\nu^2)$ solution shows how these curvature forced mid-channel bars can be a cause
650 for width variations. The dashed line in Figure 8 indicates the phase lag φ_{U_2} of the
651 symmetrical bank-excess longitudinal velocity associated with the presence of a central
652 topographical pattern. The fact that φ_{η_2} is always larger than φ_{U_2} indicates that the
653 excess near-bank longitudinal velocity locates downstream of the mid-channel bar tops.
654 This condition is almost invariably verified for values of $(\lambda, \beta, \tau_*, d_s)$ typical of gravel
655 bed rivers. As far as the near-bank excess longitudinal velocity can be assumed to be
656 positively correlated with bank erosion - as assumed in most meander migration models
657 (*Hasegawa* [1989], *Ikeda et al.* [1981]) - it can be expected that the presence of a mid-
658 channel bar would be associated with a cross-sectionally symmetrical bank erosion just a
659 few channel widths downstream the mid-channel bar top, thus determining a tendency to
660 local cross-sectional widening.

661 Analogously, a central scour hole would occur half a meander wavelength downstream
 662 the central bar deposit, triggering flow convergence which, according to model predictions,
 663 would cause a peak of the main flow thread in the central portion of the channel a few
 664 widths downstream the scour region itself. This section is expected to be subject to the
 665 maximum narrowing tendency. The overall process tends to cause planform deformation
 666 from an equiwidth meander (continuous banklines in Figure 9b) to a meander with variable
 667 width (dashed banklines) and can be referred as a laterally symmetrical mechanism for
 668 the generation of spatial width oscillations in meanders.

669 Finally, the $\mathcal{O}(\nu^2)$ solution consistently predicts an increasing relative amplitude of
 670 the mid-channel bar component of bed topography with respect to that of the point bar
 671 ($\mathcal{O}(\nu)$) for decreasing meander wavelength - i.e. increasing meander wavenumber λ (see
 672 also Figure 11a).

6.2. Effect of Width Variations on Meander Bend Growth: $\mathcal{O}(\nu\delta)$

673 The $\mathcal{O}(\nu^2)$ solution suggests that curvature variations can promote channel width vari-
 674 ations through a laterally symmetrical mechanism. The mixed $\mathcal{O}(\nu\delta)$ solution allows to
 675 investigate whether a reciprocal effect can also take place, i.e. whether the presence of
 676 width variations may affect channel curvature therefore affecting meander growth.

677 The $\mathcal{O}(\nu\delta)$ solution quantifies the correction to the linear $\mathcal{O}(\nu)$ bend stability due to
 678 the presence of width variations. This mutual nonlinear interaction gives rise to forcing
 679 terms with the same symmetry properties of the $\mathcal{O}(\nu)$ solution; laterally antisymmetric
 680 for the longitudinal momentum, water and sediments continuity equations, symmetric for
 681 the lateral momentum equation. For instance, the example term $UV_{,n}$ produces forcing

682 contributions which take the form:

$$\begin{aligned}
 V_{01}U_{10,n} &\Rightarrow (V_{1\delta}e_2 + \bar{V}_{1\delta}\bar{e}_2) (U_{1\nu,n}e_1 + \bar{U}_{1\nu,n}\bar{e}_1) \\
 &= V_{1\delta}\bar{U}_{1\nu,n}e_1 + V_{1\delta}U_{1\nu,n}e_3 + c.c.
 \end{aligned}
 \tag{16}$$

683 The above structure is reflected in the solution for \mathbf{V}_{11} . The $\mathcal{O}(\nu\delta)$ is the lowest order at
 684 which the nonlinear interaction between curvature and width forced solutions reproduces
 685 the longitudinal structure of the fundamental linear solution ($e_1 + c.c.$). This is related
 686 to the assumption $\lambda_w = 2\lambda$ that represents a key specificity of the adopted perturbation
 687 scheme. The bed topography pattern at the $\mathcal{O}(\nu\delta)$ has also a similar antisymmetric
 688 transversal pattern to that at the $\mathcal{O}(\nu)$ (Figure 12 and eqn. 12), although with different
 689 eigenfunctions, phase lags and amplitude.

In a meander with width variations, the perturbation scheme (10) indicates that
 $U(s, n = 1)$ is the sum of five main effects. The migration law (4) however suggests
 that even functions of the lateral coordinate n do not contribute to meander migration.
 Therefore only the $\mathcal{O}(\nu)$ and the $\mathcal{O}(\nu\delta)$ solutions control meander migration and stability.
 Accounting for the laterally antisymmetrical character of U_{10} and of U_{11} , it follows:

$$\zeta(s) \sim \nu [U_{10}(s, n = 1) + \delta U_{11}(s, n = 1)].
 \tag{17}$$

690 The $\mathcal{O}(\nu\delta)$ is therefore the first interaction at which an effect on bend stability can be
 691 reproduced: due to the assumption $\lambda_w = 2\lambda$, the $\mathcal{O}(\nu\delta)$ velocity perturbation has the
 692 same longitudinal structure of the fundamental $\mathcal{O}(\nu)$ perturbation.

693 The results can be summarized in a marginal stability plot in the $(\lambda - \beta)$ plane for given
 694 values of (τ_*, d_s) (Figure 10). The marginal curves in Figure 10 separate unstable (left)
 695 from stable (right) meander wavenumbers. The curve with $\delta = 0$ corresponds to the

696 result of the classical equiwidth linear bend theory - $\mathcal{O}(\nu)$ - while $\delta = 0.5$ is the upper
697 limit for geometrically meaningful planform configurations. The curves corresponding to
698 intermediate δ values indicate that, for typical aspect ratios of meandering rivers, width
699 variations tend to shift the instability region towards values of $\lambda \sim 0.2 - 0.3$, their effect
700 being stronger for intermediate values of β , between 10 and 25 for typical formative
701 conditions of gravel-bed rivers. Also the wavenumber corresponding to the maximum
702 bend amplification rate may increase. This implies that width variations are expected to
703 destabilize shorter meander bends with respect to those predicted by classical linear bend
704 theories in equiwidth meanders.

705 This may partially correct the systematical wavelength overestimation achieved by equi-
706 width linear bend theories of incipient meanders shown in Figure 10, which compares
707 predicted and measured intrinsic wavenumbers of some low sinuous rivers extracted from
708 the dataset of *Hey and Thorne* [1986]. Accounting for the presence of width variations
709 can reduce this systematic gap.

710 A second important feature of the mixed response is its tendency to produce two
711 markedly separated most unstable longitudinal modes λ associated with two distinct peaks
712 in the meander bend growth rate, represented by the function $\text{Re}(U_{11})$ (plotted in Figure
713 11b as function of meander wavelength). This behavior represents a marked difference
714 with respect to equiwidth meanders, for which bend instability is characterized by the
715 presence of a single unstable range of longitudinal modes. Meanders with intense enough
716 spatial width variations, and at sufficiently large bankfull aspect ratio, might therefore
717 display a long term tendency towards almost two different equally unstable planforms
718 characterized by well distinct meander wavelengths. *Luchi et al.* [2011] have been put

719 this in qualitative relationship with the reach-scale occurrence of chute cutoffs that has
720 been preferentially observed in meanders with spatial variations of channel width (*Lagasse*
721 *et al.* [2004] based on *Brice* [1975]).

7. Discussion

722 The implications of the outcomes of the perturbation modelling approach illustrated in
723 the previous sections are here discussed in the light of the three research questions posed
724 in the Introduction.

7.1. Role of Flow Nonlinearities in Meandering Rivers with Spatial Width Variations

725 The first research question focuses on the main physical processes that occur in rivers at
726 the meandering-braiding transition. Although the complete mathematical model accounts
727 for both linear and nonlinear processes, nonlinearity of flow and bed topography turns out
728 to be the main distinctive feature of meanders with spatial width variations, because of its
729 fundamental control on planform stability. Mid channel bars can spontaneously develop in
730 equiwidth meanders through nonlinear, curvature-forced effects and might trigger spatial
731 width oscillations. Width oscillations nonlinearly affect bend stability by shortening the
732 unstable meander bends and by determining multiple instability peaks associated with
733 two separate most unstable wavelength ranges.

734 Figure 11 combined with the field observations reported in Figure 3 can be used to show
735 that nonlinear effects are the fundamental ingredient to explain the observed differences
736 between equiwidth and wider-at-bends meanders.

737 Figure 11 couples outcomes from the nonlinear $\mathcal{O}(\nu^2)$ (a) and $\mathcal{O}(\nu\delta)$ (b) solutions.
738 Figure 11a is based on Figure 10 of *Luchi et al.* [2010b] and refers to the experimental

739 observations of *Colombini et al.* [1992] on regular sequences of small amplitude meanders
740 with constant width. It compares experimental results and model predictions for the
741 amplitude ratio between the mid-channel bar ($\mathcal{O}(\nu^2)$) and point bar ($\mathcal{O}(\nu)$) components
742 of bed topography. As verified experimentally, shorter meanders tend to have an increasing
743 dominance of the central bar component in the bed composition; they can therefore be
744 expected to exhibit a stronger tendency to develop spatial width oscillations according
745 to the autogenic, symmetrical mechanism illustrated in Figure 9b. Figure 11b (based on
746 Figure 8b of *Luchi et al.* [2010b]) shows the difference between predicted growth rate of
747 equiwidth ($\mathcal{O}(\nu)$) and wider at bends ($\mathcal{O}(\nu) + \mathcal{O}(\nu\delta)$) meandering rivers. Spatial width
748 variations cause a shift of the bend instability region towards shorter meanders, coherently
749 with the results shown in Figure 3. This can be seen from the dashed line - $\mathcal{O}(\nu\delta)$ -
750 intersecting the horizontal axis at smaller values of meander wavelength compared to
751 the continuous line - $\mathcal{O}(\nu)$. Moreover the peak in the dashed line at $L^*/W^* \sim 18 \div 20$
752 exemplifies the tendency to develop two separate most unstable wavelength ranges, which
753 is absent in linear meander models. This effect on meander wavelength selection may in
754 turn condition the nonlinear growth of meander loops with important implications also
755 for long-term simulations [*Frascati and Lanzoni, 2009*], meander loop shape [*Hooke and*
756 *Harvey, 1983*] and migration rates [*Crosato, 2008*].

757 Both the above-mentioned effects are intrinsically nonlinear and indicate that the pres-
758 ence of mid-channel bars and of spatial width variations tends to be associated with
759 shorter meander bends compared to their equiwidth counterparts. This is in good qual-
760 itative agreement with the field observations reported in Figure 3b. The presence of
761 multiple unstable longitudinal modes and of a richer content of transverse bed topogra-

762 phy modes like mid-channel bars also suggests that meanders close to the transition with
 763 braided rivers shall display a higher rate of morphological activity, which also qualitatively
 764 agrees with the observations of *Lagasse et al.* [2004] reported in Figure 3a.

765 Finally, the effect of nonlinearity in meandering morphologies close to the transition with
 766 braiding shall be accounted for into physics-based channel pattern predictors based on
 767 linear theories for free migrating [*Parker, 1976; Fredsøe, 1978*] and forced steady [*Crosato*
 768 *and Mosselman, 2009*] bars. These predictors compute the most probable lateral number
 769 m of bars for reach-averaged, representative values of discharge, sediment size, longitudinal
 770 slope and bankfull width, with the predicted value of m increasing with channel aspect
 771 ratio β . When alternate bars ($m = 1$) are the most probable, a single-thread, meandering
 772 pattern is predicted; transitional morphologies correspond to $1.5 < m < 2.5$ and braiding
 773 to $m > 3$. Nonlinearities in transitional meanders may imply the formation of mid-
 774 channel bars ($\mathcal{O}(\nu^2)$) large enough to produce a mid-channel bar configuration ($m = 2$)
 775 that corresponds to a transitional morphology for β values smaller than those predicted
 776 on the basis of linear bar theories.

7.2. Cause-Effect Relationships between Mid-Channel Bars and Width Variations in Meandering Rivers

777 Coupling linear $\mathcal{O}(\delta)$ with nonlinear $\mathcal{O}(\nu^2)$ models and examining cross sectional profiles
 778 from real meandering rivers (Figures 7 and 9a) allows to propose a possible key to dis-
 779 close the chicken-egg question about cause-effect relations between mid-channel bars and
 780 width variations in meandering rivers. A laterally symmetric and a laterally asymmetric
 781 mechanisms are proposed as possible causes for spatial width variations in curved chan-
 782 nels. Their main properties are summarized in Table 2. In the symmetrical mechanism,

783 mid-channel bars due to morphodynamic nonlinearities ($\mathcal{O}(\nu^2)$) force width variations.
784 On the contrary, width variations associated to the asymmetrical mechanism force mid
785 channel bars as a linear bed response ($\mathcal{O}(\delta)$).

786 The local unbalance between erosion and accretion processes at opposite banks can be
787 referred to as a laterally asymmetrical mechanism (Figure 7). The bank profiles in typical
788 bend sections (Figure 9a) are asymmetrical, with the eroding bank showing a nearly
789 vertical profile and the accreting bank gently sloping. Locally channel width tends to
790 change when erosion and accretion do not keep the same pace. This is largely determined
791 by bank-related processes, which are not the subject of a detail analysis in the present
792 work.

793 The symmetrical mechanism is suggested by the central bed topography pattern in the
794 cross-sectional profile shown in Figure 9b and by the lateral sidewalls being both close
795 to vertical. This suggests that opposite banks at the same cross section must be subject
796 to erosion. As sketched in Figure 9b, the steering topographical effect associated with
797 curvature-driven mid-channel bars can force the flow to diverge against both banks, thus
798 inducing laterally symmetrical bank erosion, which eventually results in cross-sectional
799 widening.

7.3. Linking Planform Stability Models with Classification of Meandering River Patterns

800 The third research question of the present work relates to the possibility of understand-
801 ing differences between meandering styles through physically-based mathematical models.
802 To make a first step in this direction, the present modelling framework can be put in qual-
803 itative relation to the form-based classification of meandering patterns originally proposed

804 by *Brice* [1975] and lately refined by *Lagasse et al.* [2004]. This is achieved in a process-
 805 oriented perspective in Figure 12, which links (arrows) the meandering styles proposed by
 806 *Lagasse et al.* [2004] (right column) with the planform stability properties as predicted at
 807 the different orders of approximation $\mathcal{O}(\nu)$, $\mathcal{O}(\delta)$, $\mathcal{O}(\nu^2)$, $\mathcal{O}(\nu\delta)$ within the present mod-
 808 elling framework. Moreover the planform and the flow-bedform pattern corresponding to
 809 each order is reported in the left column.

810 Equiwidth meandering patterns (like type-B₁) are predicted by the fundamental pro-
 811 cess of bend instability [*Ikeda et al.*, 1981], which generates Kinoshita-type meanders
 812 through geometric nonlinearities [*Seminara et al.*, 2001]. Bend instability selects long
 813 enough meander wavelengths while shorter perturbations are stable thus promoting chan-
 814 nel straightening. Multi-lobed patterns like G₁ are associated with the progressive elon-
 815 gation of meander bends which promotes the instability of higher planform harmonics
 816 (*Camporeale et al.* [2007], *Zolezzi et al.* [2009]).

817 Meandering patterns that approach the transition with braiding are characterized by
 818 the presence of spatial width variations. Following *Brice* [1975]'s classification, two major
 819 classes can be distinguished, depending on chutes being rare (classes B₂, C, E and G₂)
 820 or common (class D). The linear - $\mathcal{O}(\delta)$ - and nonlinear - $\mathcal{O}(\nu^2)$, $\mathcal{O}(\nu\delta)$ - instream pro-
 821 cesses described by the proposed modelling framework can play a relevant role in their
 822 morphodynamics.

823 The oscillations in width may be initiated due to a laterally symmetrical and a laterally
 824 antisymmetrical mechanism (Table 2). The analysis at the $\mathcal{O}(\nu^2)$ has shown that every
 825 meandering channel has an intrinsic tendency to develop mid-channel bars that eventu-
 826 ally determine symmetrical flow divergence against both banks. This is likely to trigger

827 symmetrical bank erosion slightly downstream of the mid-channel bar location, thus pro-
828 ducing longitudinal width oscillations at a double frequency with respect to curvature
829 variations. The intensity of such process is predicted to crucially depend on meander
830 wavelength, being higher for shorter meander bends.

831 On the contrary, when erosion and accretion occur at the opposite banks of the same
832 cross-section, widening (narrowing) can be produced by an asymmetrical mechanism as far
833 as the two processes do not keep the same pace [*Nanson and Hickin, 1983*]. As widening
834 and narrowing cannot indefinitely grow in time, also cyclic width variations in time can
835 be expected. Although already argued by *Parker et al. [2011]*, quantitative field evidence
836 and modelling of such temporal oscillations have not been provided so far.

837 Once initiated, longitudinal width oscillations can increase their amplitude in time
838 according to the linear planform instability mechanism detected by *Repetto et al. [2002]*,
839 who showed that only long wavelengths ($\lambda = \lambda_w/2 < 0.1$) are planimetrically unstable,
840 tend to develop more intense width oscillations possibly leading the channel to bifurcate.

841 In meandering rivers two competing autogenic mechanisms are therefore predicted to
842 compete in the temporal evolution of spatial width variations. These mechanisms act
843 oppositely for "short" and "long" meander bends. Short bends (i.e. $\lambda > 0.2 \div 0.3$, cf.
844 Figure 11) tend to have a stronger mid-channel bar component in the bed topography and
845 therefore a higher tendency to symmetrically initiate spatial width oscillations. However
846 these oscillations tend to be planimetrically stable according to the $\mathcal{O}(\delta)$ mechanism,
847 which therefore competes against their temporal growth. It must be remarked that these
848 considerations apply at the initial development stage of width oscillations and that the
849 temporal dynamics of such competition cannot be quantitatively investigated within the

850 present modelling framework. The reverse behavior is predicted for longer meander bends
851 (i.e. $\lambda < 0.1$). The geometrical properties of spatial width oscillations of B₂, C, D or E
852 meandering styles can partially reflect such opposite competition dynamics.

853 No single universal mechanism can be considered responsible for chutes initiation and
854 stability in wider at bends meandering planforms (type D in Figure 12); rather these are
855 recognized to result from either over-bar or over-bank incision, mid-channel bar growth
856 or scroll-slough development [*Grenfell et al.*, 2011]. While the dynamics of chutes can be
857 fully understood only accounting for both autogenic and allogenic processes, the role of
858 the autogenic component on at least two of the controlling processes can be related to the
859 predicted effects of width oscillations on bend instability - $\mathcal{O}(\nu\delta)$.

860 Chute initiation in large, sand-bed meandering rivers due to scroll-slough development
861 has been shown by *Grenfell et al.* [2011] to occur more frequently at bends characterized by
862 high lateral extension rates, which are predicted to increase under hydraulic conditions for
863 which spatial width variations enhance linear bend instability. Chute formation associated
864 with mid-channel bar growth [*Bridge et al.*, 1986], and possibly with over-bar incision can
865 instead be put in relation with the opposite stabilizing effect of the mixed $\mathcal{O}(\nu\delta)$ response
866 which similarly tends to reduce channel sinuosity with respect to hydraulically - equivalent
867 equiwidth meanders. Finally, and in a much broader sense, the predicted development
868 of multiple unstable longitudinal modes at the $\mathcal{O}(\nu\delta)$ can be viewed as a tendency to
869 develop two almost equally unstable meandering channels with different wavelengths,
870 whose geomorphic expression can be that of relatively stable "bifurcate" meander bends,
871 as those with chutes common.

8. Concluding Remarks

872 The present work has aimed at quantitatively investigating the autogenic morphody-
873 namic processes that shape the flow-bed topography fields in meandering channels with
874 spatial width variations, viewed as representative of river patterns at the meandering-
875 braiding transition. Three main research issues have been addressed. (1) The instream
876 physical processes associated with meandering styles close to transition with braiding; (2)
877 the cause-effect relationships associated with the contemporary presence of spatial cur-
878 vature and width variations; (3) the possibility to investigate morphological differences
879 between equiwidth and transitional meandering styles through physics-based, instream
880 morphodynamic models.

881 The analysis has been developed through (i) a review of field observations and mathe-
882 matical models relevant for meandering rivers with spatial width oscillations; (ii) a quan-
883 tification of the properties of width variations in meanders relevant for modelling purposes;
884 (iii) a unified, hierarchical derivation of models that are based on a perturbation approach
885 and that allow association of different physical processes with the mathematical solution
886 at each perturbation order; (iv) a review of the key model outcomes and comparison
887 with the few available field and laboratory observations, integrated with some original
888 model application; (v) a conceptual linkage between the model-predicted planform stabil-
889 ity properties and descriptive classification of meandering channel patterns.

890 The depth-averaged model is solved through a comprehensive two-parameters pertur-
891 bation approach. Results indicate that curvature-width interactions are fundamentally
892 nonlinear: the hierarchical model structure allows to disentangle cause-effect relation-
893 ships between width variations and mid-channel bars at the first nonlinear interaction.

894 In meanders with initially constant width, curvature nonlinearly forces mid-channel bar
895 growth, promoting symmetrical bank erosion further downstream, thus being a possible
896 trigger for width oscillations. These in turn can significantly affect meander stability and
897 growth and therefore condition the curvature dynamics. Nonlinearities cause wider-at-
898 bends meandering rivers to develop significantly shorter bends and to be morphologically
899 more active with respect to constant width meanders. This picture is coherent with the
900 relatively few available field observations.

901 The analysis is restricted to "autogenic" or "instream" processes, i.e. those occurring in
902 the central flow region, where the boundary effects of banks can be ignored; bank processes
903 are therefore only indirectly accounted for through a linear relationship between the near-
904 bank excess velocity and the lateral channel migration rate. Despite this simplification,
905 decades of meander modelling research have shown the potential of models relying on this
906 assumption to capture relevant physical processes. It is unquestionable, however, that the
907 full interaction dynamics between the river channel and the entire river corridor needs to
908 be understood and modelled in order to build a more complete river-floodplain dynamics.
909 This sets the scene for future research directions.

910 The array of biophysical processes and mutual interactions that need to be addressed can
911 be conveniently illustrated moving from the central flow region to the vegetated floodplain
912 patches, along an ideal cross-section of the whole river corridor. Among these, we deem
913 that the following key ingredients deserve priority in future research.

914 1. The present review, namely in Sections 2 and 3, has highlighted the relative paucity
915 of quantitative field data on spatial width variations in the various styles of meandering
916 rivers, particularly close to the transition with braiding. Spatial and temporal width

917 variations can be quantified by fully exploiting the potential of remote sensing technologies
918 and possible correlations with flow and sediment regime, catchment, floodplain and soil
919 properties shall be assessed.

920 2. While flow dynamics in the central flow region has been extensively investigated,
921 near-bank flow processes are still partially known, both for the eroding and for the ac-
922 creting bank. Their knowledge is required for full coupling of bed-banks evolution. The
923 growing scientific attention [*Güneralp et al.*, 2012] is presently biased towards processes
924 near the eroding bank (*Kean and Smith* [2006], *Blanckaert and de Vriend* [2005]).

925 3. Evidence has been provided [*Lauer and Parker*, 2008] of local sediment imbalances
926 in meandering rivers in association with bank erosion and channel curvature, with other
927 field observations [*Gautier et al.*, 2007] indicating their morphodynamic relevance. Lat-
928 eral sediment exchanges between the central flow region and the evolving banks have
929 seldom been accounted for in meandering river models and potentially have fundamental
930 implications for spatial and temporal width oscillations [*Chen and Duan*, 2006]. Mor-
931 phodynamic coupling of banks and central flow regions through analytical models also
932 show a bias towards the eroding bank, sometimes with great physical detail on processes
933 causing collapse, failure and fluvial erosion [*Motta et al.*, 2012]. The modelling framework
934 of *Parker et al.* [2011], accounts for the sediment imbalance through a laterally-integrated
935 formulation of sediment continuity at both bank regions, thus allowing the eroding bank
936 and depositing bank talk to each other and adjust the width of the central flow region
937 in space and time. Coupling with analytical morphodynamic models for the central flow
938 region [*Imran et al.*, 1999] has been only very recently attempted by *Eke and Parker*

939 [2011], coherently with the need to adopt a nonlinear model for the central flow region as
940 suggested by the present review.

941 4. The eroding-bank-bias can be attributed to the largely more advanced understand-
942 ing of bank erosion [*Rinaldi and Darby*, 2008] with respect to accretion [*Crosato*, 2008].
943 Consistent improvements in understanding and modelling the dynamics of bank accretion
944 can be expected by accounting for flow unsteadiness and sediment heterogeneity [*Brau-*
945 *drick et al.*, 2009], which have seldom been incorporated in meander migration models
946 (but see *Camporeale and Ridolfi* [2010]). The other key element controlling bank accre-
947 tion are biotic-abiotic interactions associated with riparian vegetation dynamics [*Perucca*
948 *et al.*, 2007], which plays a fundamental morphodynamic role in different types of river
949 systems [*Gurnell et al.*, 2012] and is clearly recognized to made a difference in establishing
950 single-thread and anabranching river channels on Earth [*Davies and Gibling*, 2011].

951 5. Descriptive classifications and rational predictors of channel pattern [*Kleinhans and*
952 *van den Berg*, 2011] suggest the close association of chute cutoffs with meandering river
953 styles at the transition with braiding (type D in Figure 12). Only recently significant
954 progress towards quantitative understanding of their dynamics has been achieved from
955 detailed analysis of field data (*Constantine et al.* [2009], *Micheli and Larsen* [2010], *Gren-*
956 *fell et al.* [2011]), while improvements in morphodynamic modelling of these "bifurcate
957 meander bends" are still awaited. Coupled with field-based process understanding, this
958 can decisively contribute to detect key differences along with more hidden similarities in
959 processes typical of single- and multiple-thread alluvial streams.

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Figure 1. Illustrative examples of river morphologies across the continuum of alluvial channel patterns. (a) Meandering, Rio Yurua (Amazon region). Meandering with considerable spatial width variations: (b) Rio Beni (Bolivia); (c) Sacramento River (California, U.S.). (d) Braiding: Tagliamento River (Italy). Meandering patterns with width variations, mid-channel bars and chutes can be viewed as transitional forms between single- and multiple-thread patterns.

Figure 2. Modified Brice classification of single-thread alluvial river patterns. Classes have been grouped according to the absence or presence of spatial variations in channel width. Adapted from Lagasse et al. (2004).

Figure 3. Cumulative percentage of apex bend movement (a) and of dimensionless meander wavelength L^*/W^* (b) separately plotted for equiwidth (class B_1) and wider-at-bends (class C) meandering river bends of Brice classification. Based on Lagasse et al. (2004) and Luchi et al.(2011).

Table 1. Classification of existing mathematical models able to predict physical

processes occurring in meandering rivers with spatial width variations.

	Model category	Reference	Time and spatial scales of application	Key issues in relation to paper subject
Analytical Models	(1) Instream morphodynamics in straight and meandering channels with spatial width variations and Partheniades-type lateral migration law	<i>Repetto et al.</i> [2002]; <i>Luchi et al.</i> [2010b]; <i>Luchi et al.</i> [2011]; <i>Luchi et al.</i> [2012]; <i>Frascati and Lanzoni</i> [2011]	Reach scale and wavelength scale, taken as representative of a longer reach with homogeneous hydro-morphological properties; flow-bed at equilibrium with the given planform	Explicitly address the separate role of the two planform forcings and their nonlinear interactions
	(2) Meander migration models with physics-based bank erosion migration and analytical flow-bed topography field	<i>Chen and Duan</i> [2006] <i>Motta et al.</i> [2012] <i>Parker et al.</i> [2011] <i>Eke and Parker</i> [2011]	Reach scale and wavelength scale, taken as representative of a longer reach with homogeneous hydro-morphological properties; flow-bed at equilibrium with the given planform; can simulate different relative timescales of bed and banks evolution	Focus on meander migration; emphasis on bank processes, mostly on erosion, with varying degrees of detail. Some incorporate lateral sediment exchange associated with banklines shift
Numerical Models	(3) Morphodynamic models with physics-based bank erosion and numerical flow-bed topography field	<i>Mosselman</i> [1992]; <i>Darby et al.</i> [2002]; <i>Jang and Shimizu</i> [2005]; <i>Rüther and Olsen</i> [2007];	Event-based or multiple formative events; reach scale (from one to some meander wavelengths). Suitable to simulate short-term planform evolution	Highly simplified description of bank accretion; difficult to systematically assess the separate roles of channel curvature and width variations in dependence of reach-averaged hydraulic conditions
	(4) Reduced - complexity models	<i>Coulthard and Wiel</i> [2006]; <i>Jagers</i> [2003]	Potentially suitable for long-term evolution of entire floodplain reaches because of smaller computational effort compared to fluid-mechanic based numerical models	Cellular and object based models. Potential for comparison with channel pattern prediction; physics is strongly simplified compared to other categories

Figure 4. (a) Image of a reach of the Rio Beni (Bolivia). (b) Channel width variations in the same reach. The width peaks almost invariably once per bend, regardless of bend orientation, suggesting that $L^* = 2L_w^*$. Capital letters (A, \dots, J) indicate inflection points.

Figure 5. Box plot illustrating the dimensionless amplitude δ of width variations (a) and ν of curvature variations (b) for C-class bends reported in the dataset of *Lagasse et al.* [2004]. The upper and lower limits of the grey boxes coincide with the 84th and 16th percentile respectively, with vertical lines extending to the 95th and to the 5th percentiles of the category. The percentage of river bends within the whole examined dataset for each class of λ is reported in brackets below the horizontal axis.

Figure 6. Sketch of a meandering channel with spatially variable channel width and notations employed in the mathematical model.

Figure 7. Typical asymmetric cross section in a meander bend, with opposite banks experiencing different processes, eventually resulting in local widening or narrowing. (data from River Bollin, UK; from Luchi et al., 2010)

Figure 8. (a) Aerial view of the reach of the River Bollin (NW England) used for the application of the $\mathcal{O}(\nu^2)$ solution; the yellow bar indicates the cross section plotted in Figure 9b. (b,d) Downstream views of the same inflection region in April 2008 and July 2009 (courtesy of J. Hooke). (c) Relative position of the mid channel bar top $\eta_{20}^{(2)}(n=0)$ and of the peak of the related symmetrical perturbation of the longitudinal velocity $U_{20}^{(2)}$ at the banks ($\beta = 6.21, \tau_* = 0.06, d_s = 0.03$). The yellow bar indicates the position of the cross section where the mid-channel bar top has been observed along half meander wavelength.

Figure 9. (a) Representative cross sections at meander inflection in the reach of the River Bollin investigated by Luchi et al. (2009c). (b) Conceptual explanation of the topographically-driven mechanism for the generation of width variations in a meandering channel.

Figure 10. (a) Marginal bend stability curves in the $\lambda - \beta$ parameter space for different intensities δ of spatial width variations ($\tau_* = 0.1, d_s = 0.08, \omega = 0$). (b) Comparison between the wavenumber selected by the classical linear bend theory ($\mathcal{O}(\nu)$, see Section 4.2) and the wavenumber of low sinuosity stream reaches extracted from the dataset of Hey and Thorne (1986).

Figure 11. (a) Ratio of the amplitudes of the mid-channel bar ($\mathcal{O}(\nu^2)$) and point bar ($\mathcal{O}(\nu)$) components of bed topography in the experimental runs of *Colombini et al.* [1991] (open circles) and predicted by the reviewed models for the same hydraulic and geometric conditions (closed circles). Input parameters ranges: β ($11 \div 22$); τ_* ($0.04 \div 0.07$); d_s ($0.05 \div 0.09$); based on *Luchi et al.* [2010b]. (b) Predicted meander bend growth rate for equiwidth ($\mathcal{O}(\nu)$, continuous line) and wider at bends ($\mathcal{O}(\nu) + \mathcal{O}(\nu\delta)$, dashed line) $\tau_* = 0.08$; $d_s = 0.04$; $\beta = 20$; $\delta = 0.17$. Based on *Luchi et al.* [2011].

<i>Laterally asymmetrical mechanism</i>	<i>Laterally symmetrical mechanism</i>
one bank in erosion, the opposite bank in accretion	both banks in erosion
width variations have a <i>forcing</i> function with respect to mid-channel bars	width variations have a <i>following</i> function with respect to mid-channel bars
mid-channel bars are a <i>linear</i> topographical response	mid-channel bars are a <i>nonlinear</i> topographical response

Table 2. Comparison between the two mechanisms proposed to explain the occurrence of spatial width variations in meandering channels.

Figure 12. A process-form diagram of width - curvature interactions in equiwidth and transitional meanders based on a two-parameters perturbation approach. The left column qualitatively illustrates the planform, flow and bed topography pattern corresponding to each perturbation order. The Brice meander pattern classification, partially redrawn in the right column, is put in relationship with the planform stability properties of each perturbation order.